

Thermodynamics of the fermion gas

a. The grand canonical partition function is

$$\begin{aligned}\Xi &= \prod_r (1 + e^{-\beta(\epsilon_r - \mu)}) \\ \ln \Xi &= \sum_r \ln(1 + e^{-\beta(\epsilon_r - \mu)}) = \int_0^\infty d\mathcal{E} G(\mathcal{E}) \ln(1 + e^{-\beta(\mathcal{E} - \mu)}),\end{aligned}$$

but

$$\Pi = -pV = -k_B T \ln \Xi \quad \text{so} \quad p(T, \mu)V = k_B T \int_0^\infty G(\mathcal{E}) \ln(1 + e^{-\beta\mathcal{E}} e^{\beta\mu}) d\mathcal{E}.$$

b. For spin- $\frac{1}{2}$ particles, $G(\mathcal{E}) = V \left[\frac{2m^3}{\pi^4 \hbar^6} \right]^{1/2} \sqrt{\mathcal{E}}$. Thus

$$\begin{aligned}p(T, \mu) &= k_B T \left[\frac{2m^3}{\pi^4 \hbar^6} \right]^{1/2} \int_0^\infty \sqrt{\mathcal{E}} \ln(1 + e^{-\beta\mathcal{E}} e^{\beta\mu}) d\mathcal{E} \quad [\dots \text{use the substitution } x = \beta\mathcal{E} \dots] \\ &= (k_B T)^{5/2} \left[\frac{2m^3}{\pi^4 \hbar^6} \right]^{1/2} \int_0^\infty \sqrt{x} \ln(1 + e^{-x} e^{\beta\mu}) dx.\end{aligned}$$

c. To evaluate

$$\int_0^\infty \sqrt{x} \ln(1 + e^{-x} e^{\beta\mu}) dx$$

integrate by parts using

$$\begin{aligned}du &= \sqrt{x} dx & v &= \ln(1 + e^{-x} e^{\beta\mu}) \\ u &= \frac{2}{3} x^{3/2} & dv &= \frac{dx}{1 + e^{-x} e^{\beta\mu}} (-e^{-x} e^{\beta\mu}) = -\frac{dx}{e^x e^{-\beta\mu} + 1}\end{aligned}$$

and find

$$\int_0^\infty \sqrt{x} \ln(1 + e^{-x} e^{\beta\mu}) dx = \left[\frac{2}{3} x^{3/2} \ln(1 + e^{-x} e^{\beta\mu}) \right]_0^\infty + \frac{2}{3} \int_0^\infty \frac{x^{3/2}}{e^x e^{-\beta\mu} + 1} dx.$$

As $x \rightarrow \infty$, $x^{3/2} \ln(1 + e^{-x} e^{\beta\mu}) \rightarrow x^{3/2} e^{-x} e^{\beta\mu} \rightarrow 0$, so the term in square brackets vanishes and we have

$$p(T, \mu) = \frac{2}{3} (k_B T)^{5/2} \left[\frac{2m^3}{\pi^4 \hbar^6} \right]^{1/2} \int_0^\infty \frac{x^{3/2}}{e^x e^{-\beta\mu} + 1} dx.$$

d.

$$\begin{aligned}E(T, V, \mu) &= \int_0^\infty G(\mathcal{E}) f(\mathcal{E}) \mathcal{E} d\mathcal{E} \\ &= V \left[\frac{2m^3}{\pi^4 \hbar^6} \right]^{1/2} \int_0^\infty \sqrt{\mathcal{E}} \frac{1}{e^{\beta(\mathcal{E} - \mu)} + 1} \mathcal{E} d\mathcal{E} \\ &= V (k_B T)^{5/2} \left[\frac{2m^3}{\pi^4 \hbar^6} \right]^{1/2} \int_0^\infty \frac{x^{3/2}}{e^x e^{-\beta\mu} + 1} dx\end{aligned}$$