

Quantum Mechanics

Model Solutions for Sample Exam for Second Examination

1. According to equation (5.35) in the text,

$$|\psi(t)\rangle = e^{-(i/\hbar)Et} \left[e^{-(i/\hbar)(-A)t} \langle e_1 | \psi(0) \rangle |e_1\rangle + e^{-(i/\hbar)(+A)t} \langle e_2 | \psi(0) \rangle |e_2\rangle \right].$$

where $|\psi(0)\rangle = |d\rangle$. But from (5.34), $\langle e_1 | d \rangle = -\frac{1}{\sqrt{2}}e^{+i\phi}$ and $\langle e_2 | d \rangle = +\frac{1}{\sqrt{2}}e^{+i\phi}$ so

$$|\psi(t)\rangle = e^{-(i/\hbar)Et} \left(\frac{1}{\sqrt{2}}e^{+i\phi} \right) \left[-e^{-(i/\hbar)(-A)t} |e_1\rangle + e^{-(i/\hbar)(+A)t} |e_2\rangle \right].$$

We must first find the amplitude

$$\begin{aligned} \langle d | \psi(t) \rangle &= e^{-(i/\hbar)Et} \left(\frac{1}{\sqrt{2}}e^{+i\phi} \right) \left[-e^{-(i/\hbar)(-A)t} \langle d | e_1 \rangle + e^{-(i/\hbar)(+A)t} \langle d | e_2 \rangle \right] \\ &= e^{-(i/\hbar)Et} \left(\frac{1}{\sqrt{2}}e^{+i\phi} \right) \left(\frac{1}{\sqrt{2}}e^{-i\phi} \right) \left[-e^{-(i/\hbar)(-A)t}(-1) + e^{-(i/\hbar)(+A)t}(+1) \right] \\ &= e^{-(i/\hbar)Et} \left(\frac{1}{2} \right) \left[e^{-(i/\hbar)(-A)t} + e^{-(i/\hbar)(+A)t} \right] \\ &= e^{-(i/\hbar)Et} \cos \left(\frac{A}{\hbar}t \right). \end{aligned}$$

Thus the probability of finding the nitrogen atom “down” is

$$\cos^2 \left(\frac{A}{\hbar}t \right).$$

2. From the commutator,

$$\hat{x}\hat{p} - \hat{p}\hat{x} = i\hbar \quad \text{so} \quad \langle \hat{x}\hat{p} \rangle - \langle \hat{p}\hat{x} \rangle = i\hbar.$$

But

$$(\hat{x}\hat{p})^\dagger = \hat{p}^\dagger \hat{x}^\dagger = \hat{p}\hat{x} \quad \text{so} \quad \langle \hat{x}\hat{p} \rangle^* = \langle \hat{p}\hat{x} \rangle.$$

Together

$$\langle \hat{x}\hat{p} \rangle - \langle \hat{x}\hat{p} \rangle^* = i\hbar \quad \text{so} \quad 2i\Im\{\langle \hat{x}\hat{p} \rangle\} = i\hbar \quad \text{so} \quad \Im\{\langle \hat{x}\hat{p} \rangle\} = \frac{1}{2}\hbar.$$

3. We have

$$\hat{A}\hat{B} - \hat{B}\hat{A} = c\hat{B} \quad \text{so} \quad \hat{A}\hat{B}|a\rangle - a\hat{B}|a\rangle = c\hat{B}|a\rangle \quad \text{so} \quad \hat{A}(\hat{B}|a\rangle) = (a+c)(\hat{B}|a\rangle).$$

That last equation says that $\hat{B}|a\rangle$ is an eigenvector of \hat{A} with eigenvalue $a+c$. Hence it is also an eigenvector of \hat{A}^3 with eigenvalue $(a+c)^3$.

4. Lorentzian wavepacket

For

$$\psi(x) = \frac{A}{x^2 + \gamma^2} e^{ikx},$$

the normalization is given through

$$\begin{aligned} 1 &= \int_{-\infty}^{+\infty} \psi^*(x)\psi(x) dx \\ &= A^2 \int_{-\infty}^{+\infty} \frac{1}{(x^2 + \gamma^2)^2} dx \quad \text{[[Use Dwight 120.2 giving...]]} \\ &= A^2 \left[\frac{x}{2\gamma^2(x^2 + \gamma^2)} + \frac{1}{2\gamma^3} \arctan \frac{x}{\gamma} \right]_{-\infty}^{+\infty} \\ &= A^2 \left[\frac{1}{2\gamma^3} \left(\frac{\pi}{2} - \left(-\frac{\pi}{2} \right) \right) \right] = A^2 \frac{\pi}{2\gamma^3} \end{aligned}$$

so $A^2 = 2\gamma^3/\pi$.

The mean kinetic energy is given through

$$\langle \text{KE} \rangle = -\frac{\hbar^2}{2m} \int_{-\infty}^{+\infty} \psi^*(x) \frac{d^2}{dx^2} \psi(x) dx = +\frac{\hbar^2}{2m} \int_{-\infty}^{+\infty} \frac{d\psi^*(x)}{dx} \frac{d\psi(x)}{dx} dx$$

where we have used integration by parts to cast the integral into a more symmetric form. (And to avoid taking a second derivative!) Now

$$\begin{aligned} \frac{d\psi(x)}{dx} &= -\frac{A 2x}{(x^2 + \gamma^2)^2} e^{ikx} + ik\psi(x) \\ \frac{d\psi^*(x)}{dx} &= -\frac{A 2x}{(x^2 + \gamma^2)^2} e^{-ikx} - ik\psi^*(x) \end{aligned}$$

so

$$\begin{aligned} \langle \text{KE} \rangle &= \frac{\hbar^2}{2m} \int_{-\infty}^{+\infty} \left[\frac{A 2x}{(x^2 + \gamma^2)^2} e^{-ikx} + ik\psi^*(x) \right] \left[\frac{A 2x}{(x^2 + \gamma^2)^2} e^{ikx} - ik\psi(x) \right] dx \\ &= \frac{\hbar^2}{2m} \left[\int_{-\infty}^{+\infty} \left(\frac{A 2x}{(x^2 + \gamma^2)^2} \right)^2 dx + \int_{-\infty}^{+\infty} \left(\frac{A 2x}{(x^2 + \gamma^2)^2} e^{-ikx} \right) \left(-ik \frac{A}{x^2 + \gamma^2} e^{ikx} \right) dx \right. \\ &\quad \left. + \int_{-\infty}^{+\infty} \left(+ik \frac{A}{x^2 + \gamma^2} e^{-ikx} \right) \left(\frac{A 2x}{(x^2 + \gamma^2)^2} e^{ikx} \right) dx + k^2 \int_{-\infty}^{+\infty} \psi^*(x)\psi(x) dx \right] \end{aligned}$$

Of these four integrals, the last one is just the normalization integral, so it is one. The second and third integrals have odd integrands, so they are zero. We're left with

$$\langle \text{KE} \rangle = \frac{\hbar^2}{2m} \left[4A^2 \int_{-\infty}^{+\infty} \frac{x^2}{(x^2 + \gamma^2)^4} dx + k^2 \right].$$

Now, using Dwight 122.4,

$$\int_{-\infty}^{+\infty} \frac{x^2}{(x^2 + \gamma^2)^4} dx = \left[-\frac{x}{6(x^2 + \gamma^2)^3} + \frac{x}{24\gamma^2(x^2 + \gamma^2)^2} + \frac{x}{16\gamma^4(x^2 + \gamma^2)} + \frac{1}{16\gamma^5} \arctan \frac{x}{\gamma} \right]_{-\infty}^{+\infty} = \frac{\pi}{16\gamma^5},$$

whence

$$\langle \text{KE} \rangle = \frac{\hbar^2}{2m} \left[k^2 + \frac{1}{2\gamma^2} \right].$$